

# 9 Modalanalyse, Eigenvektor

Begriffsbildung:  $\bar{\Psi}$  - Eigenvektor;  $\bar{\Psi} = [\psi_1; \dots; \psi_n]^T$   
 $\bar{\Psi}_n$  - n-te Eigenform;  $\bar{\Psi}_n = [\psi_{1n}, \psi_{2n}, \dots, \psi_{nn}]^T$

Bestimmung des Eigenvektors

aus Ansatz  $(A - \lambda E) \bar{\Psi} \cdot e^{j\omega t} = \bar{0}$

$(A - \lambda E) \bar{\Psi} = \bar{0}$  Eigenvektorgleichung

Bsp. f. 2  $\lambda_1$  und  $\lambda_2$  sind berechnet

① für  $\lambda_1$   $(A - \lambda_1 E) \bar{\Psi}_1 = \bar{0}$  (1)

② für  $\lambda_2$   $(A - \lambda_2 E) \bar{\Psi}_2 = \bar{0}$  (1)

(1)  $(A - \lambda_1 E) \begin{bmatrix} \psi_{11} \\ \psi_{21} \end{bmatrix} = \bar{0}$

(2)  $(A - \lambda_2 E) \begin{bmatrix} \psi_{12} \\ \psi_{22} \end{bmatrix} = \bar{0}$

(1)  $\begin{bmatrix} \frac{c}{m} - \lambda_1 & -\frac{c}{2m} \\ -\frac{c}{m} & \frac{c}{m} - \lambda_1 \end{bmatrix} \begin{bmatrix} \psi_{11} \\ \psi_{21} \end{bmatrix} = \bar{0}$  für für  $\lambda_1$

$\left. \begin{aligned} (\frac{c}{m} - \lambda_1) \psi_{11} - \frac{c}{2m} \psi_{21} &= 0 \\ -\frac{c}{m} \psi_{11} + (\frac{c}{m} - \lambda_1) \psi_{21} &= 0 \end{aligned} \right\} \psi_{21} = \sqrt{2} \cdot \psi_{11} \text{ linear abhängig!}$

z.B.:  $\psi_{11} = 1 \Rightarrow \psi_{21} = \sqrt{2} \quad \bar{\Psi}_1 = k_1 \begin{bmatrix} 1 \\ \sqrt{2} \end{bmatrix} = k_1 \cdot \underline{\underline{\bar{\Psi}_1^N}}$

(2)  $\begin{bmatrix} \frac{c}{m} - \lambda_2 & -\frac{c}{2m} \\ -\frac{c}{m} & \frac{c}{m} - \lambda_2 \end{bmatrix} \begin{bmatrix} \psi_{12} \\ \psi_{22} \end{bmatrix} = \bar{0}$

$\left. \begin{aligned} (\frac{c}{m} - \lambda_2) \psi_{12} - \frac{c}{2m} \psi_{22} &= 0 \\ -\frac{c}{m} \psi_{12} + (\frac{c}{m} - \lambda_2) \psi_{22} &= 0 \end{aligned} \right\} \psi_{22} = -\sqrt{2} \cdot \psi_{12} \text{ linear abhängig!}$

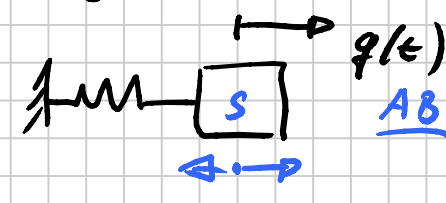
z.B.  $\psi_{12} = 1 \Rightarrow \psi_{22} = -\sqrt{2} \quad \bar{\Psi}_2 = k_2 \begin{bmatrix} 1 \\ -\sqrt{2} \end{bmatrix} = k_2 \cdot \underline{\underline{\bar{\Psi}_2^N}}$

? Warum sind  $\bar{\Psi}_1$  und  $\bar{\Psi}_2$  beliebig normierbar?

Ansatz:  $\bar{q} = \bar{\gamma} e^{j\omega_0 t} = \bar{\gamma}_1 e^{j\omega_{01} t} + \bar{\gamma}_2 e^{j\omega_{02} t} + \dots + \bar{\gamma}_n e^{j\omega_{0n} t}$

physikalische Interpretation

$\bar{q} \hat{=}$  Wegvektor  
 $\bar{\gamma} \hat{=}$  Wegvektoren



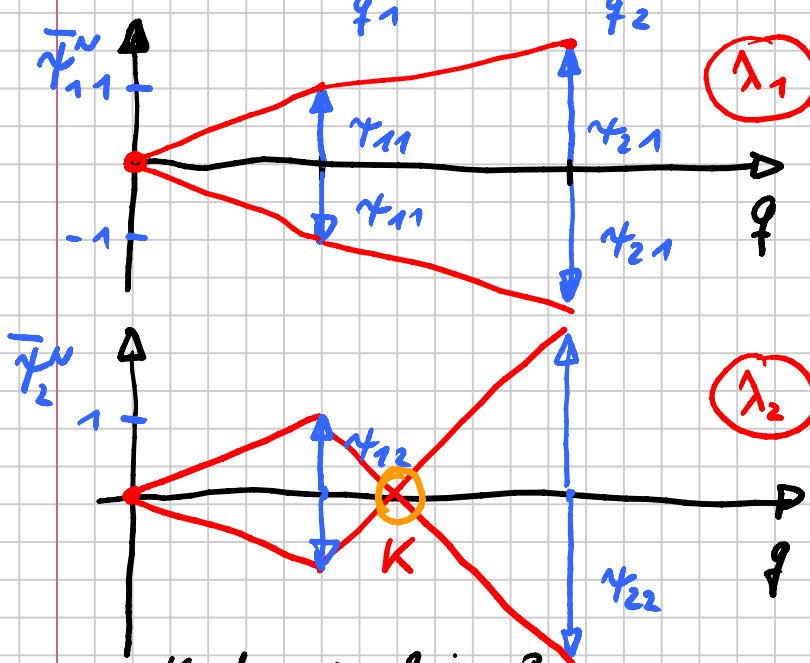
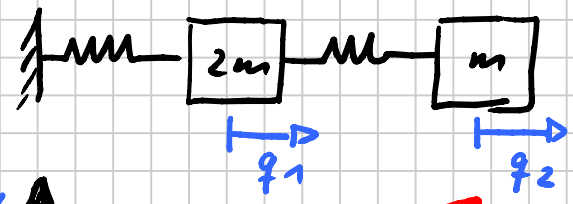
AB:  $q(t=0) = q_0$   
 $\dot{q}(t=0) = v_0$

Lösung  $\bar{q}(t)$  unter AB

$\bar{q}(t) = \bar{C}_1 e^{j\omega_{01} t} + \dots + \bar{C}_n e^{j\omega_{0n} t}$   
 $\bar{C}_1 = k_1 \bar{\gamma}_1^N ; \bar{C}_n = k_n \bar{\gamma}_n^N$

grafische Darstellung der Eigenformen

Bsp. f=2



$\gamma_{11} = 1 ; \gamma_{21} = \sqrt{2}$   
 $\gamma_{11} = -1 ; \gamma_{21} = -\sqrt{2}$

Knotenbild

$\gamma_{12} = 1 ; \gamma_{22} = -\sqrt{2}$   
 $\gamma_{12} = -1 ; \gamma_{22} = \sqrt{2}$

Knoten K: keine Bewegung

0 Knoten Grundwelle  $\leadsto \lambda_1 \leadsto \omega_{01}$  Grundwelle  
 1 Knoten 1. Oberwelle  $\leadsto \lambda_2 \leadsto \omega_{02}$  1. Oberwelle  
 Prüfen!

$\bar{q} = \sum_n \bar{C}_n e^{j\omega_{0n} t}$   $\neq$  nicht das Knotenbild

